

# Rectennas Revisited

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**Abstract**—In the late 1960s, a new concept was proposed for an infrared absorbing device called a “rectenna” that, combining an antenna and a nanoscale metal–insulator–metal diode rectifier, collects electromagnetic radiation in the terahertz regime, with applications as detectors and energy harvesters. Previous theories hold that the diode rectifies the induced terahertz currents. Our results, however, demonstrate that the Seebeck thermal effect is the actual dominant rectifying mechanism. This new realization that the underlying mechanism is thermal-based, rather than tunneling-based, can open the way to important new developments in the field, since the fabrication process of rectennas based on the Seebeck effect is far simpler than existing processes that require delicate tunnel junctions. We demonstrate for the first time the fabrication of a rectenna array using an efficient parallel transfer printing process featuring nearly one million elements.

**Index Terms**—Antenna arrays, infrared (IR) detectors, rectifiers.

## I. INTRODUCTION

NUMEROUS groups [1]–[34] since the sixties have been working in the field of antennas coupled to metal–insulator–metal (MIM) diodes, namely rectennas, for absorbing infrared (IR) radiation. The MIM diode is a quantum device in which a nanometer-thick dielectric, often an oxide, is sandwiched between two electrodes featuring dissimilar work functions. The dissimilarity in the work functions is thought to cause an asymmetric tunnel current to flow through the oxide with respect to the polarity of the electrodes. According to the

existing theory, the MIM diode is coupled in parallel to the antenna, so the high-frequency terahertz currents induced in the antenna under illumination are partially rectified by the MIM diodes. Theory holds that using dissimilar work functions for the electrode metals enables rectenna operation without the need for an externally applied bias, which would be the case for energy harvesting. Also, for rectennas as IR detectors, the antenna response may be increased by biasing the MIM diode with a constant external voltage, in which case the two electrodes may have the same work function.

Currently, several groups are working on theory and experiments to design the “ultimate” rectenna suitable for industrial applications, such as energy harvesting and IR detection [26]–[29], [31], [34]. This includes the topic of impedance matching in order to extract electrical currents efficiently without having high losses [33]. In this paper, we use two complementary approaches, namely an antenna array fabricated by a novel nanotransfer-printing process, discussed below, and individual antennas fabricated by electron-beam lithography and lift-off techniques to demonstrate that the Seebeck effect is the dominant rectification mechanism in rectennas, in contradiction to the established view that attributes signal response to the rectification of MIM diode currents. By eliminating the need for the tunnel barrier, our finding leads to a significant simplification of the design and fabrication of rectennas and, therefore, brings a major breakthrough in rectenna applications for IR detectors, solar cells and energy harvesting devices.

Moreover, we present the first rectenna array featuring almost 1 million dipole antennas, coupled to nanometer-size MIM junctions, which we will show function due to thermal rather than tunneling effects. We report here the response of several sub-assemblies of antenna-coupled MIM diodes to incident IR radiation. Further, the optoelectronic characteristics of single antennas and arrays of antennas were simulated [35] and tested to minimize destructive interference of neighboring antennas. Simulations showed that the signal strength for 10.6  $\mu\text{m}$  radiation was a maximum for antennas spaced about 5  $\mu\text{m}$  apart. By fabricating our rectenna arrays with our highly efficient nanotransfer-printing (nTP) process, we demonstrate that our rectenna arrays can be manufactured at low-cost on virtually any surface.

This paper reports in Section II on the material properties with respect to the MIM diode, and on the antenna array design. The electronic and optoelectronic properties of rectenna arrays are discussed in Section III. In Section IV, the focus is placed on the Seebeck effect. Biased rectennas are discussed in Section IV. A final conclusion is given in the last section (see Section V).

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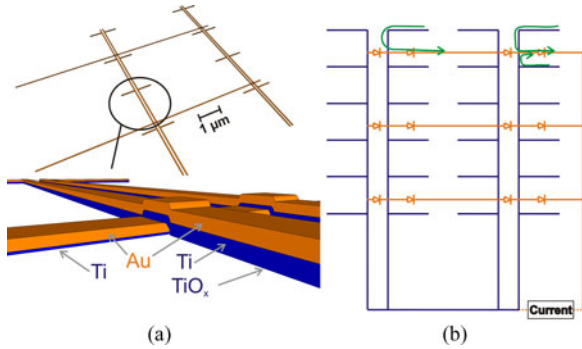


Fig. 1. (a) MIM diodes are formed at the intersections between the Au bottom electrodes (with the Ti adhesion promoter beneath) and the Ti top electrodes (with the 2.2-nm-thick plasma-grown  $\text{TiO}_x$  dielectric beneath and the Au delamination promoter on top). Each MIM diode is coupled to an antenna (Ti-Au). (b) The resulting circuit diagram shows the lines transferred in the first nTP step (horizontal lines), the antenna structures transferred in the second step (vertical lines including antenna structures) and the path taken by the rectified currents (indicated by the pointers).

## II. MATERIAL PROPERTIES AND ANTENNA ARRAY DESIGN

In our previous investigations of the MIM diode phenomenon, we fabricated aluminum (Al) [36], titanium (Ti) [37], gold (Au) [36], and platinum (Pt) [32] as electrode materials and native grown aluminum oxide ( $\text{AlO}_x$ ) [38], plasma-grown  $\text{AlO}_x$  [36], plasma grown titanium oxide ( $\text{TiO}_x$ ) [37], and self-assembled monolayers (SAMs) as inorganic or organic dielectric materials. Our previous studies [39] showed that organic materials or native oxides as a sole insulator does not form a uniform dielectric layer. At that time, we sought the highest nonlinearities and lowest resistance in the current–voltage ( $I$ – $V$ ) characteristics of the diodes, as well as the lowest permittivity of the dielectric. During the investigation of these materials [37], we found that the dielectric permittivity is strongly dependent on the dielectric thickness [36], [37]. The most promising results for overall diode characteristics were obtained with plasma-grown  $\text{AlO}_x$  (3.6 nm thick) and plasma-grown  $\text{TiO}_x$  (2.2 nm thick) [37]. Based on our investigations, we chose Ti– $\text{TiO}_x$ –Au as the optimum material combination for obtaining nonlinear MIM diode properties.

Previously, we fabricated only individual rectennas [30]. However, in order to consider a device for energy harvesting applications, an array of rectennas must be demonstrated. The relative distances between the antennas, as well as the dimensions of the individual antennas, were simulated and optimized for sensitivity to a wavelength of  $10.6 \mu\text{m}$  and to avoid destructive interferences between neighboring antenna elements (see Fig. 1).

A two-step nTP process was designed to efficiently fabricate an array of almost one million dipole antennas, all coupled to MIM junctions. We printed the devices using templates, or “stamps,” directly onto a transparent glass substrate that can be easily activated by an oxygen plasma treatment and thus provides a hydrophilic surface. It is important to note that our fabrication process consists of efficient, industry-compatible steps, and does not require repeated long and expensive exposure times. The first stamp, shown in Fig. 2(a), transfers Au

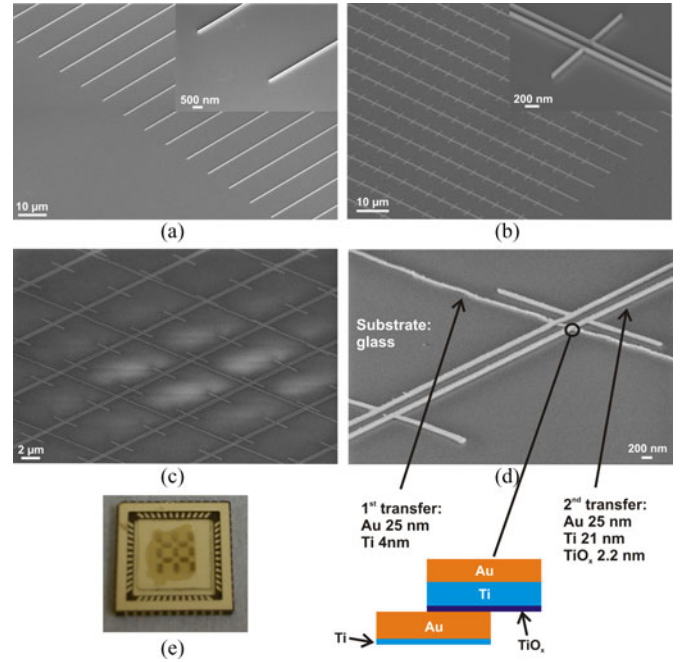


Fig. 2. (a) Scanning electron microscopy (SEM) images of the first stamp and (b) of the second stamp. (c) Magnified view of the rectenna array transferred onto a glass substrate. (d) High-resolution image of the rectenna array, showing the antenna structures and the MIM junctions created at the overlap between the Au bottom electrodes and the Ti top electrodes. (e) Packaged chip including the contact pads fabricated by optical lithography.

bottom electrodes of the MIM diodes in the form of narrow lines onto the glass substrate, with a thin Ti layer that promotes adhesion to the glass surface. The metal antennas and the  $\text{TiO}_x$ –Ti stacks that form the tunnel dielectric and the top electrodes are then transferred from a second stamp [see Fig. 2(b)], creating antenna-coupled Au– $\text{TiO}_x$ –Ti junctions in the overlap areas. The Ti top electrode is covered by a thin Au layer that is needed to facilitate delamination of the  $\text{TiO}_x$ –Ti stack from the stamp. Fig. 2(c) shows a portion of the antenna array, and Fig. 2(d) shows a single antenna including the MIM junction. This low-cost and high-speed fabrication technique offers a critical advantage over conventional rectenna fabrication by electron-beam lithography. The rectenna arrays fabricated with nTP have a surface that consists only of Au, and thus features excellent properties with respect to the absorption of irradiated radiation [40]. After the fabrication of the array, we define contact pads on top of the array by conventional optical lithography and lift-off techniques in order to characterize individual sectors featuring 25 000 rectennas each. A full chip in a package is shown in Fig. 2(e).

## III. ELECTRONIC AND OPTOELECTRONIC PROPERTIES OF THE RECTENNA ARRAY

In order to determine the rectenna response of several sectors of the array, we connected individual sectors (as defined by the bonding pads) to a voltage amplifier, and illuminated the entire array with a linearly polarized  $\text{CO}_2$  laser operating at  $10.6 \mu\text{m}$ , i.e., 28 THz. A half-wave plate was used to adjust the orientation of the linear polarization of the incident radiation. We observed

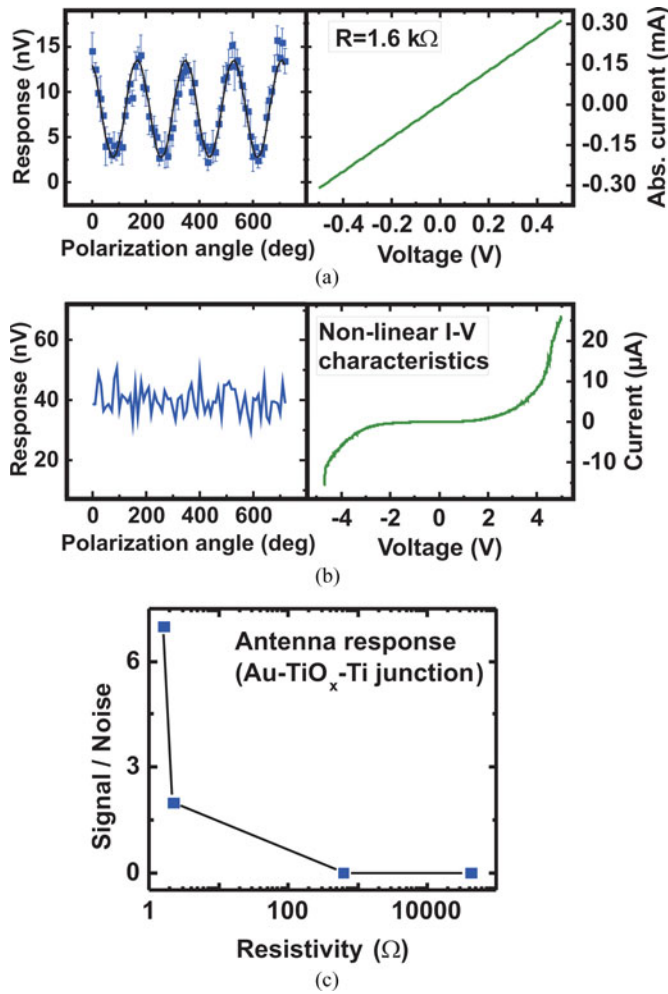


Fig. 3. (a) Polarization-dependent antenna response of a section of the rectenna array that shows linear  $I$ - $V$  characteristics with a small resistance. (b) In contrast, the antenna response of a section of the rectenna array that shows high-resistance and exponential current-voltage characteristics is below the noise level. (c) The highest signal to noise ratio is attained by a sector with lowest resistivity, which supports the model of a thermal rectification effect in the rectenna array.

a clear polarization-dependent antenna response in several sectors. When the linear polarization of the incident electric field was perpendicular to the axis of the dipole antennas, the antenna response was near the measureable noise level. When the polarization of the IR wave was parallel to the axis of the antennas, the measured response was seven times larger than the noise level [see Fig. 3(a)]. To our knowledge, an antenna array of this size has never before been demonstrated to produce a signal at THz frequencies. In addition, by optimizing the design of the rectenna array by simulations, we were able to obtain the highest polarization ratio for unbiased rectennas in this wavelength regime. Polarization ratio is defined as the ratio of the antenna response to incident radiation with linear polarization parallel to the antenna axis to that of polarization perpendicular to the antenna axis.

However, other sectors investigated did not exhibit any radiation response. In order to understand this result, we investigated the  $I$ - $V$  characteristics of the diode ensembles without illumina-

tion. In those sectors where no antenna response was observed, the expected tunnel-diode-like, non-linear  $I$ - $V$  behavior was measured, and had a resistance in the range of megaohms. In contrast, we found that the sectors in which we observed a clear antenna response exhibited a linear  $I$ - $V$  characteristic with a small resistance ( $\sim 1 \text{ k}\Omega$ ), that is, the dielectric in several sectors exhibited linear resistive paths rather than tunnel-barrier behavior. These linear resistive paths might occur due to either a defect in the oxide or a mechanical break in the oxide. We believe that the higher noise value in the high-resistance sectors is due to Johnson noise because the Johnson noise is proportional to the resistance [41], [42]. In other words, the observed relationship between the  $I$ - $V$  characteristics of the MIM junctions and the measured antenna response is exactly opposite to that expected by the original MIM diode theory. Even when considering surface plasmon polaritons to be present due to IR absorption at the metal-dielectric interface, the tunnel current is not high enough in order to detect a dc signal [43], [44]. This suggests that the operation mechanism of the rectennas is not based on tunneling rectification in the MIM junctions, as was previously thought.

#### IV. SEEBECK EFFECT VERSUS TUNNELING RECTIFICATION

We believe that the Seebeck effect [45] is the dominant mechanism in our devices that do not exhibit a tunnel barrier, leading to the antenna response. Due to IR absorption, the array is heated up, whereas the electrodes furthest away from the array and closest to the contact pads are not heated. The temperature gradient in each of the electrodes creates a voltage drop, which is related to the Seebeck coefficient of the electrode metals [46]. Further, the Seebeck-effect provides a consistent explanation for the observed relationship between the current-voltage behavior of the sectors of the antenna array and the measured antenna response (quantified by the signal-to-noise ratio) seen in Fig. 3. Devices with a lower resistivity exhibit a larger signal-to-noise ratio [see Fig. 3(c)]. After the fabrication of the device, we characterized parts of the array using a CO<sub>2</sub> laser beam setup. The device under test (DUT) was mounted on a translation stage that can shift the DUT perpendicular to the incident laser beam in order to scan the whole sample. The CO<sub>2</sub> laser is illuminating an area of about  $1 \text{ mm}^2$  and the laser is operated at its main emission line at  $10.6 \mu\text{m}$  (28 THz). A polarizer ensures a linear polarization of the emitted light and a half-wave plate can rotate the polarization by  $360^\circ$ . A chopper is modulating the laser beam and the response of the DUT is determined with a lock-in amplifier that is synchronized to the modulation frequency.

In order to verify our findings, we fabricated test structures of single rectennas consisting of various noble metals in direct contact with one another, i.e., without a tunnel barrier. Since noble metals do not form a native oxide, no insulating layer is present in these rectennas [see Fig. 4(a)]. All antennas were characterized with respect to their optoelectronic (antenna response under illumination) and electronic ( $I$ - $V$  characteristics) properties. We compared the antenna response of several rectennas featuring the same dimensions but different metal combinations [see Fig. 4(a)], and found a linear dependence of the antenna response on the difference in the Seebeck coefficients of the two

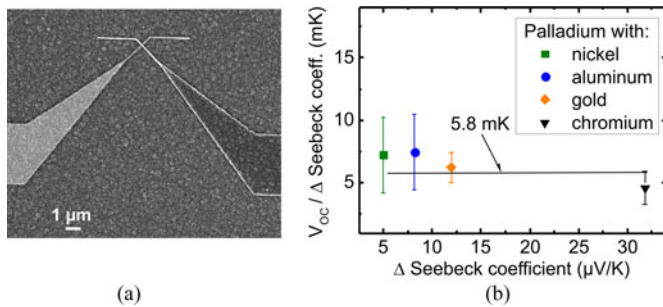


Fig. 4. (a) SEM image of single rectenna consisting of only noble metals without oxide layer. (b) A constant ratio between the open-circuit voltage (VOC) and relative Seebeck coefficients is obtained, and is equal to the average temperature difference between the junctions.

electrodes [see Fig. 4(b)], in excellent agreement with theory:

$$V = (S_{M1} - S_{M2}) \times (T_{\text{hot}} - T_{\text{cold}})$$

where  $S_{M1}$  and  $S_{M2}$  are the Seebeck coefficients of the two electrode metals and  $T_{\text{hot}}$  and  $T_{\text{cold}}$  are the temperatures of the metals at the hot and the cold junction. The precise Seebeck coefficients are known from previous investigations in which we included thermometer structures close to the individual rectenna structures. Since we did not incorporate thermometer structures into the present antenna array, we infer from Fig. 4(b) that the average temperature difference between the two junctions is the same for all of the tests, and is approximately 5.8 mK. Invoking the rectification mechanism of rectennas with the Seebeck effect also explains why the response of the functioning rectenna sectors is small. Since the difference in the Seebeck coefficients of Au and Ti is relatively small, the induced open-circuit voltage is consequently also smaller than for other possible material combinations with a larger difference of their Seebeck coefficients. Also, the array was not designed to maximize temperature differences across the MIM junctions. If rectennas are built based on the Seebeck effect without the need for a tunnel barrier, the fabrication process is drastically simplified, since the manufacture and implementation of the ultra-thin oxide layers represented one of the key challenges for these devices.

## V. ANTENNA RESPONSE IN BIASED RECTENNAS

Believing that the tunneling mechanism is the dominant rectifying mechanism, several groups proposed to bias the rectenna in order to increase the asymmetry of the tunneling currents around the bias voltage, and thus, to increase the rectified current [15], [19] according to the dependence of rectified current on the nonlinearity of the  $I$ - $V$  operating point,

$$I_{\text{dc}} = \frac{1}{4} \frac{d^2 I}{dV^2} \Big|_{V=V_{\text{bias}}} V_0^2$$

where  $V_0$  is the amplitude of the induced voltage and  $V_{\text{bias}}$  is the applied voltage.

We also followed this principle and biased ensembles of antennas showing nonlinear  $I$ - $V$  characteristics and a lack of response to incident IR radiation at zero bias. When applying an external bias, we obtained an antenna response [see Fig. 5(a)].

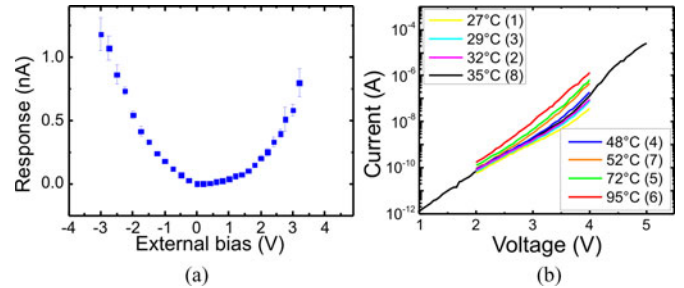


Fig. 5. (a) When externally biasing a sector of the rectenna array that showed no antenna response without an applied bias, an antenna response in the form of a current signal is observed. (b) The temperature-dependent current-voltage characteristics of a MIM diode show that the resistance of the MIM junction decreases with increasing temperature, which is the signature of a bolometer effect.

However, further investigations showed that the response obtained under external bias can also be explained by a standard bolometric effect [47] due to the temperature dependence of the tunnel current. When radiation is incident on the biased rectennas, a resistance change at the MIM junction occurs due to heating. This explanation is supported by our temperature dependent  $I$ - $V$  measurements of our MIM junctions [see Fig. 5(b)]. Therefore, the present concept of diode rectification for biased rectennas should be further investigated in light of the contribution due to temperature dependent tunneling.

## VI. CONCLUSION

We have reported on the rectification mechanism in IR-sensitive rectennas. Our investigations show that rectification of the terahertz currents induced by the absorbed incident IR radiation is not due to tunneling rectification in the unbiased MIM tunnel junctions, but rather due to the Seebeck effect that results from the temperature gradient in the two different metals of our rectenna structure. We suggest that the response to incident radiation of biased antennas coupled to MIM junctions can also be attributed to thermal effects. Future investigations should give further basis to this statement and show that the ac-signal conversion into dc is more efficient due to thermal effects than to tunnel effects. We further showed a highly efficient method for the fabrication of nanoscale rectenna arrays of nearly one million elements. Our findings offer a new concept that should be implemented into new designs of IR rectenna detectors and solar rectenna arrays for energy harvesting.

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